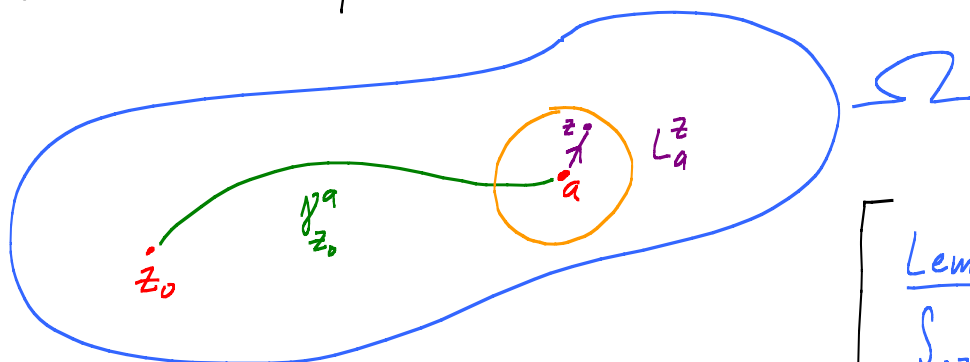


Lesson 13 The Cauchy Theorem

Assuming $\int_{\gamma} f dz = 0$ for all closed γ in a domain Ω (f analytic)
we want to cook up F with $F' = f$



$$\begin{aligned} F(z) &= \int_{\gamma_z} f(w) dw \\ &= \left(\int_{\gamma_a^{z_0}} + \int_{L_a^z} \right) f dw \\ &\quad \uparrow \\ &\quad \text{Constant} \end{aligned}$$

Lemma:

$$\begin{aligned} \int_{\gamma_A^z} c dw &= c(z-A) \\ \parallel \\ \int_{\gamma_A^z} \frac{d}{dw} [cw] dw &= cz - cA \checkmark \end{aligned}$$

Want: $F'(z) = \frac{d}{dz} \left(\int_{L_a^z} f dw \right)$

f analytic $\Rightarrow f$ continuous. So near a :

$$f(z) = f(a) + E(z) \quad \text{where } E(z) \rightarrow 0 \text{ as } z \rightarrow a.$$

$$\begin{aligned} DQ &= \frac{F(z) - F(a)}{z - a} = \frac{\int_{L_a^z} f dw}{z - a} \\ &= \frac{\int_{L_a^z} (f(a) + E(w)) dw}{z - a} \\ &= \underbrace{\frac{f(a)(z-a)}{z-a}}_{f(a)} + \underbrace{\frac{1}{z-a} \int_{L_a^z} E(w) dw}_{E(z)} \end{aligned}$$

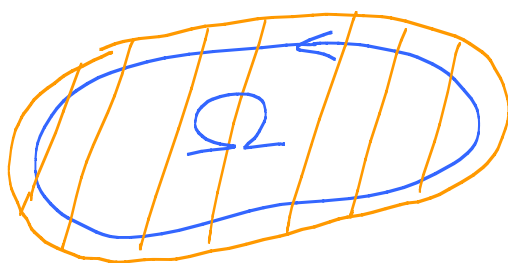
where $|E(z)| = \frac{1}{|z-a|} \left| \int_{L_a^z} E(w) dw \right|$

$$\leq \frac{1}{|z-a|} \cdot \underbrace{\left(\max_{L_a^z} |E(w)| \right)}_{\rightarrow 0 \text{ as } z \rightarrow a} \underbrace{\text{Length}(L_a^z)}_{|z-a|}$$

So $E(z) \rightarrow 0$ as $z \rightarrow a$ and we get $F'(a) = f(a)$.

Read 4.4 b, skip a.

Cauchy's Theorem: Suppose f is "analytic inside and on" a simple closed curve γ . Then $\int_{\gamma} f dz = 0$.



γ

f analytic on 
 $\Omega = \text{inside of } \gamma$

Green's Theorem: $\int_{\gamma} P dx + Q dy = \iint_{\Omega} \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dx dy$

Pf of Cauchy's Thm:

$$f(x+iy) = u(x,y) + i v(x,y)$$

$$z(t) = x(t) + i y(t)$$

$$\int_{\gamma} f dz = \int_a^b f(z(t)) z'(t) dt$$

$$= \int_a^b \left[u(x(t), y(t)) + i v(x(t), y(t)) \right] (x'(t) + i y'(t)) dt$$

$$= \int_a^b (u \dot{x} - \overbrace{v \dot{y}}^{i^2 = -1}) dt + i \int (u \dot{y} + v \dot{x}) dt$$

$$= \int_{\gamma} u dx - v dy + i \int_{\gamma} u dy + v dx$$

Green's

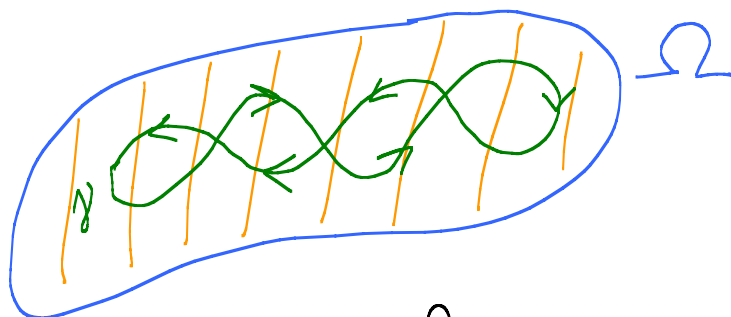
Note: $\int_a^b u(x(t), y(t)) \underbrace{\frac{dx}{dt} dt}_{dx} = \int_{\gamma} u dx$

$$= \iint_{\Omega} \underbrace{\left(\frac{\partial}{\partial x} (-v) - \frac{\partial u}{\partial y} \right)}_{u_y = -v_x} dx dy + i \iint_{\Omega} \underbrace{\left(\frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} \right)}_{u_x = v_y} dx dy$$

$$= 0$$

Cauchy's Thm #2: Suppose Ω is a domain with no holes (Ω is "simply connected"), and f analytic on Ω . Then $\int_{\gamma} f dz = 0$ for every closed γ in Ω .

Why:



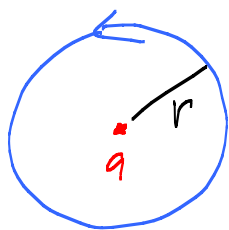
$$\int_{\gamma} f dz = \text{Sum of } \int_{\text{loops}} f dz = 0 \text{ by Cauchy's \#1.}$$

Cor: Using "Big fact" about getting F with $F' = f$

$\Leftrightarrow \int_{\gamma} f dz = 0$ for closed γ , we see that

analytic fns on a domain with no holes have analytic antiderivatives.

Important calculation:



$$C_r(a) : z(t) = a + re^{it} \\ 0 \leq t \leq 2\pi$$

$$\int_{C_r(a)} \frac{1}{z-a} dz = \int_0^{2\pi} \frac{1}{(a+re^{it})-a} \underbrace{ire^{it}}_{z'(t)} dt$$

$$= i \int_0^{2\pi} 1 dt = 2\pi i$$

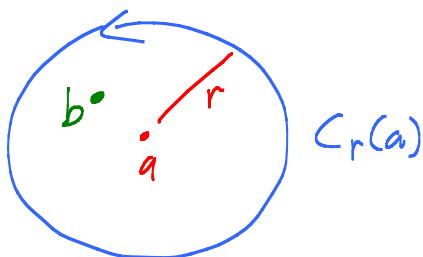
Consequence: Cauchy Thm fails on $\mathbb{C} - \{a\}$ $\leftarrow$ domain with a hole

Also $\frac{1}{z-a}$ has no analytic antiderivative on $\mathbb{C} - \{a\}$.

$$\left[\int_{C_r(a)} F' dz = F(\text{END}) - F(\text{START}) = 0 \right]$$

Of course: $\frac{d}{dz} \log_a(z-a) = \frac{1}{z-a}$ $\leftarrow$ fails to hook up on $\mathbb{C} - \{a\}$.

Big fact:



$$\int_{C_r(a)} \frac{1}{z-b} dz = \begin{cases} 2\pi i & b \text{ inside } C_r(a) \\ 0 & b \text{ outside} \end{cases}$$

$\uparrow$ by Cauchy

Why: Let $F(w) = \int_{C_r(a)} \frac{1}{z-w} dz$

$$DQ = \int DQ \rightarrow \int_{C_r(a)} \frac{d}{dw} \left[\frac{1}{z-w} \right] dz$$

$$\begin{aligned} F'(w) &= \int_{C_r(a)} \frac{1}{(z-w)^2} dz \\ &= \int_{C_r(a)} \frac{d}{dz} \left[\frac{-1}{z-w} \right] dz \\ &= \bigcirc \end{aligned}$$

Aha! $F' \equiv 0$. So F is constant.

Since $F(a) = 2\pi i$, the constant is $2\pi i$ ✓

Remarks: Stokes' $\Rightarrow$ Green's

$$\int_{\gamma} \omega = \iint_{\Omega} d\omega$$

$\omega = Pdx + Qdy$ ← differential 1-form

$$\begin{aligned} d\omega &= dP \wedge dx + dQ \wedge dy \\ &= \left(\frac{\partial P}{\partial x} dx + \frac{\partial P}{\partial y} dy \right) \wedge dx + \left(\frac{\partial Q}{\partial x} dx + \frac{\partial Q}{\partial y} dy \right) \wedge dy \\ &= \frac{\partial P}{\partial x} \underbrace{dx \wedge dx}_{=0} + \frac{\partial P}{\partial y} \underbrace{dy \wedge dx}_{=-dx \wedge dy} + \frac{\partial Q}{\partial x} dx \wedge dy + \frac{\partial Q}{\partial y} \underbrace{dy \wedge dy}_{=0} \\ &= \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) \underbrace{dx \wedge dy}_{\text{element of area } dA} \end{aligned}$$