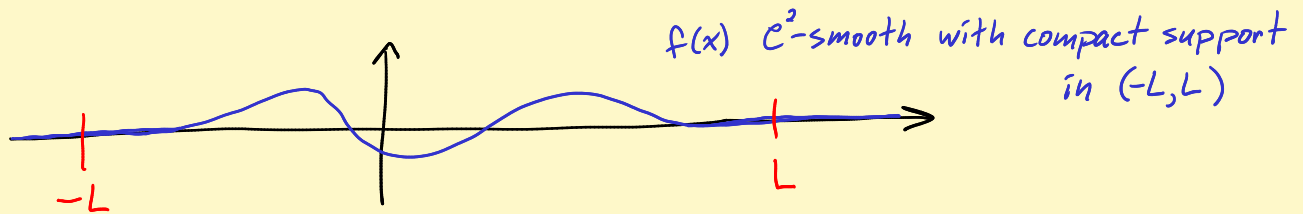


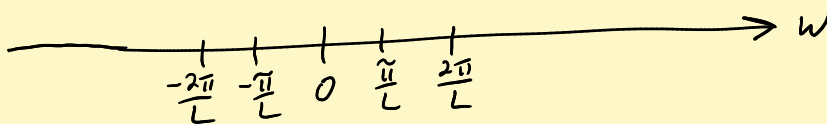
# Lecture 28 The complex version of the Fourier transform



$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{in\pi x/L} \quad \text{where} \quad c_n = \frac{1}{2L} \int_{-L}^L f(t) e^{-in\pi t/L} dt$$

$$= \frac{1}{\Delta w} \sum_{n=-\infty}^{\infty} \left( \frac{1}{2L} \int_{-L}^L f(t) e^{-i \frac{n\pi}{L} t} dt \right) e^{i \frac{n\pi}{L} x}$$

$\frac{\pi}{L} = \Delta w$



$$f(x) = \xrightarrow{L \rightarrow \infty} \frac{1}{2\pi} \int_{-\infty}^{\infty} \left( \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt \right) e^{i\omega x} d\omega$$

where to put?

Fourier transform

## Conventions

$$\hat{f}(s) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-isx} dx$$

also write  $\hat{f}(s) = \mathcal{F}[f]$

Fourier inversion formula  $f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(s) e^{isx} ds$

Also write  $\check{g}(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(s) e^{isx} ds$

$$= \mathcal{F}^{-1}[g]$$

Inversion formula :  $f = \check{\check{f}} = \mathcal{F}^{-1}[\mathcal{F}[f]]$

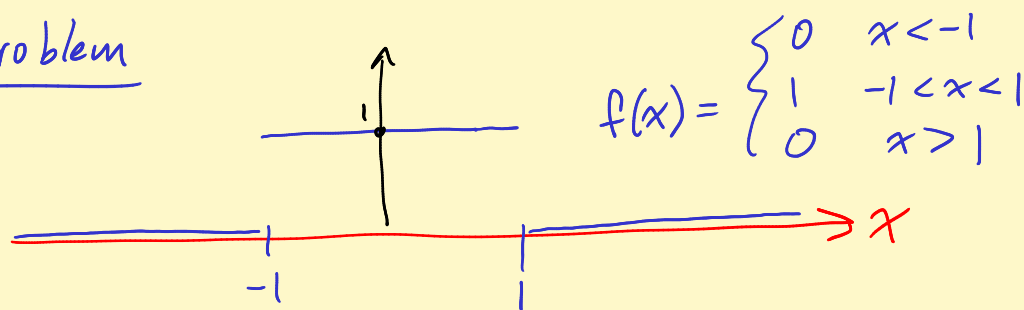
Big properties 1)  $\mathcal{F}[f'] = is \hat{f}(s)$

2)  $\mathcal{F}[f''] = -s^2 \hat{f}(s)$

3) If  $f$  is piecewise  $C^1$ -smooth and  $\int_{-\infty}^{\infty} |f| dx < \infty$ ,

then  $f(x) \stackrel{""}{=} \mathcal{F}^{-1}[\hat{f}] = \begin{cases} f(x) & f \text{ cont at } x \\ \text{midpoint of jump at jumps} \end{cases}$

Hot wire problem



$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2} \quad (c=1)$$

$$u(x, 0) = f(x)$$

(No boundary. So no boundary conditions.)

Big idea: Apply  $\mathcal{F}$  in space variable  $x$ .

Write  $\hat{u}(s, t) = \mathcal{F}[u] = \frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} u(x, t) e^{-isx} dx$

Hmmm. Take  $\frac{\partial}{\partial t}$  of this  $\uparrow$

$$\frac{\partial \hat{u}(s, t)}{\partial t} = \frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} \underbrace{\frac{\partial u}{\partial t}(x, t)}_{\frac{\partial^2 u}{\partial x^2}} e^{-isx} dx$$

$$= \mathcal{F}\left[\frac{\partial^2 u}{\partial x^2}\right] = -s^2 \mathcal{F}[u]$$

Get

$$\boxed{\frac{\partial \hat{u}}{\partial t}(s, t) = -s^2 \hat{u}(s, t)}$$

← An "ODE" with a parameter  $s$  floating along

$$\hat{u}(s, t) = \underbrace{C(s)}_{\text{constant might depend on } s} e^{-s^2 t}$$

constant might depend on  $s$ .

Also, ok for  $C(s)$  to be complex #.

Next step: Get  $C(s)$  from initial condition.

$$\underbrace{\hat{u}(s, 0)}_{= C(s)} = \frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} \underbrace{u(x, 0)}_{f(x)} e^{-isx} dx$$

Aha!

$$\begin{aligned}
 C(s) = \hat{f}(s) &= \frac{1}{\sqrt{2\pi i}} \int_{-1}^1 1 \cdot e^{-isx} dx \\
 &= \frac{1}{\sqrt{2\pi i}} \left[ \frac{1}{(-is)} e^{-isx} \right]_{x=-1}^1 \\
 &= \frac{1}{\sqrt{2\pi i}} \left[ \frac{e^{-is}}{-is} - \frac{e^{is}}{(-is)} \right] \\
 &= \frac{1 \cdot 2}{\sqrt{2\pi i}} \left[ \frac{e^{is} - e^{-is}}{2i} \right] \cdot \frac{1}{s} \\
 &= \sqrt{\frac{2}{\pi i}} \frac{\sin s}{s}
 \end{aligned}$$

So  $\hat{u}(s, t) = \sqrt{\frac{2}{\pi i}} \frac{\sin s}{s} e^{-s^2 t}$

Last step: Undo  $\hat{\cdot}$ . Get

$$\begin{aligned}
 u(x, t) &= \frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} \left( \sqrt{\frac{2}{\pi i}} \frac{\sin s}{s} e^{-s^2 t} \right) e^{isx} ds \\
 &= \frac{1}{\pi i} \int_{-\infty}^{\infty} \frac{\sin s}{s} e^{-s^2 t} \left( \cos sx + i \sin sx \right) ds
 \end{aligned}$$

↑  
even
↑  
even
↑  
even
↑  
odd

Symmetric interval :  $\int_{-\infty}^{\infty} = \lim_{A \rightarrow \infty} \int_{-A}^A$

(even)(even)(odd) = odd      So imaginary part of  $\int_{-\infty}^{\infty} = 0$ .

sol<sup>n</sup>

$$u(x,t) = \frac{1}{\sqrt{11}} \int_{-\infty}^{\infty} \frac{\sin s}{s} e^{-s^2 t} \cos sx \, ds$$

$$= \frac{2}{\sqrt{11}} \int_0^{\infty} \frac{\sin s}{s} e^{-s^2 t} \cos sx \, ds$$



As soon as  $t > 0$ , solution becomes  $C^\infty$ -smooth and integrals converge rapidly. Heat dissipates quickly and  $u(x,t) \rightarrow 0$  uniformly in  $x$  as  $t \rightarrow \infty$ .