

Lesson 11 Solution to the Dirichlet problem on the disc

Formula sheet on Homepage

Parseval's

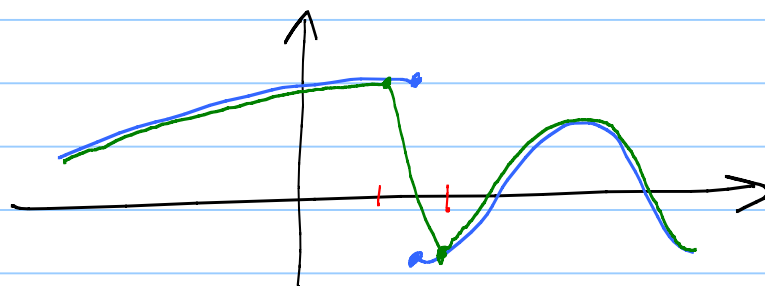
$$0 \leq \underbrace{\left\| f - \sum_{-N}^N c_n e^{inx} \right\|^2}_{\|f\|^2 - \frac{1}{2\pi} \sum_{-N}^N |c_n|^2} \leq \underbrace{\left\| f - \sum_{-N}^N A_n e^{inx} \right\|^2}_{\sigma_N(x)} \quad \text{can make } < \varepsilon \text{ by Fejér's.}$$

Last time: Assumed f cont. 2π -periodic

$$\text{Fejér's} \Rightarrow \text{Parseval's} \quad \|f\|^2 = \frac{1}{2\pi} \sum_{-\infty}^{\infty} |c_n|^2.$$

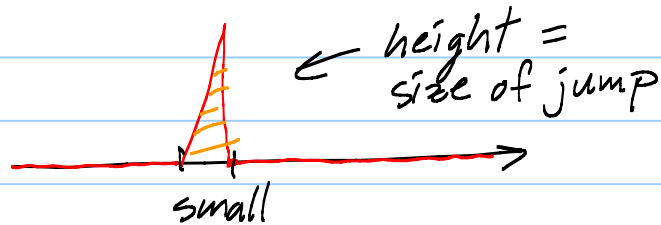
Hmmm: Bessel's holds when f is piecewise cont.

Trick



$\| \text{Green} - \text{Blue} \|^2$ can be made small!

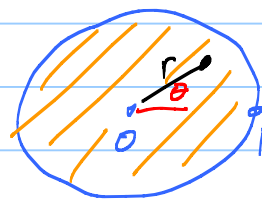
$$\int_{-\pi}^{\pi} | \text{Green} - \text{Blue} |^2 dx :$$



small area!

Conclusion: Parseval's holds for piecewise continuous f .

Solve heat problem on a metal disc:



$D_1(0)$

Know temp on boundary (circle).

Temp = $u(r, \theta)$. Let it reach "steady state".

Math problem: Need 1) $\Delta u = 0$ on $D_1(0)$

2) u continuous up to circle,

3) $u(1, \theta) = f(\theta)$ $0 \leq \theta < 2\pi$

$$\begin{cases} f(0) = f(2\pi) \\ f'(0) = f'(2\pi) \end{cases}$$

$$\Delta u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2}$$

Defⁿ: $D_1(0) = \{ (x, y) : \sqrt{x^2 + y^2} < 1 \} = \{ (r, \theta) : r < 1 \}$

Try $u(r, \theta) = R(r)\Theta(\theta)$

$$\Delta u = 0$$

$$R''\Theta + \frac{1}{r}R'\Theta + \frac{1}{r^2}R\Theta'' \equiv 0$$

divide by $R\Theta$. Multiply by r^2 .

$$\boxed{\frac{r^2 R'' + r R'}{R} = -\frac{\Theta''}{\Theta}} = \lambda$$

Both sides must = same constant λ . [Let $\theta = \frac{\pi}{2}$. Let $r = \frac{1}{2}$.]

Two ODEs:
$$\begin{cases} \Theta'' + \lambda \Theta = 0 \\ r^2 R'' + r R' - \lambda R = 0 \end{cases}$$

(H)-problem: $(H)'' + \lambda (H) = 0$

periodic boundary conditions: $\begin{cases} (H)(0) = (H)(2\pi) \\ (H)'(0) = (H)'(2\pi) \end{cases}$

$u(r, \theta) = R(r) (H)(\theta) \leftarrow \text{no jumps, no kinks at } 0 \sim 2\pi.$

$\lambda = 0$: $(H)''(\theta) \equiv 0$

$(H)(\theta) = A\theta + B$

$(H)(0) = B \overset{\text{want}}{=} A \cdot 2\pi + B$
see $\boxed{A=0.}$

So $(H)(\theta) \equiv B$

Oh good! $(H)'(0) = 0 = (H)'(2\pi)$

Get constant solⁿ: $(H)(\theta) \equiv B$. Take $B=1$.

Write $\boxed{(H)_0(\theta) = 1}$ for $\lambda = 0$.

$\lambda < 0$: Get $(H)(\theta) = c_1 e^{k\theta} + c_2 e^{-k\theta}$ $\lambda = -k^2$.

Ouch! Can't satisfy periodic B.C.

Get no solⁿs when $\lambda < 0$.

$\lambda > 0$: Get periodic solⁿs exactly when $\lambda = n^2$
 $n=1, 2, 3, \dots$

$(H)_n(\theta) = a_n \cos n\theta + b_n \sin n\theta$

Back to $R(r)$ problem: $r^2 R'' + r R' - \lambda R = 0$

Euler eqn!

Case $\lambda = 0$: $r^2 R'' + r R' = 0$. Let $V = R'$

(Standard form) (*) $\frac{dV}{dr} + \frac{1}{r} V = 0$ First order linear!

Int. factor $\mu = e^{\int \frac{1}{r} dr} = e^{\ln r} = r$

Mult (*) by μ : $r \frac{dV}{dr} + 1 \cdot V = 0$

$\frac{d}{dr} [\underbrace{\mu V}_{rV}]$

Get $rV = c_1$, a const.

$V = \frac{c_1}{r}$

So $R' = \frac{c_1}{r}$

$R = c_1 \underbrace{\ln r}_{\text{nasty term}} + c_2$

Physical considerations:
 $\Rightarrow c_1 = 0$.

Get $R(r) \equiv c_2$. Take $c_2 = 1$.

Also get $u_0(r, \theta) = R_0(r) \Theta_0(\theta) = 1$

Case $\lambda = n^2$: $r^2 R'' + r R' - n^2 R = 0$

Euler: Try $R(r) = r^p$

$R'(r) = p r^{p-1}$

$R''(r) = p(p-1) r^{p-2}$

Plug in

$$r^2 (p(p-1) r^{p-2}) + r (p r^{p-1}) - n^2 r^p \stackrel{\text{want}}{=} 0$$

$$\left[\underbrace{p(p-1) + p - n^2}_{\text{need } = 0} \right] r^p = 0$$

$$p^2 - n^2 = 0$$

$$\boxed{p = \pm n}$$

Get $R(r) = c_1 \underbrace{r^{-n}}_{\text{nasty!}} + c_2 r^n$
Must have $c_1 = 0$.

$$R_n(r) = r^n$$

$$u_n(r, \theta) = R_n(r) \Theta(\theta) = r^n (a_n \cos n\theta + b_n \sin n\theta)$$