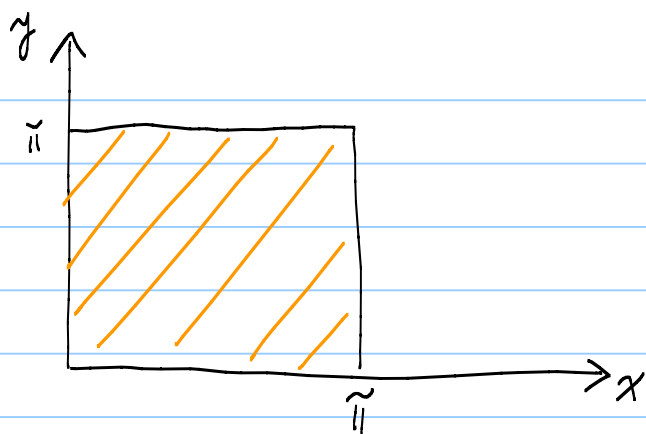


Lesson 25 Vibrating square drum, Laplace transform

HWK 5 due Wed



$$\frac{\partial^2 u}{\partial t^2} = c^2 \Delta u \quad (c=1)$$

$$u(x, y, t) = 0 \text{ on frame}$$

$$\begin{cases} u(x, y, 0) = f(x, y) \\ \frac{\partial u}{\partial t}(x, y, 0) = g(x, y) \end{cases}$$

$$u(x, y, t) = \sum_{n, m=1}^{\infty} \sin nx \sin my \left(A_{nm} \cos \omega_{nm} t + B_{nm} \sin \omega_{nm} t \right)$$

$$\text{where } \omega_{nm} = \sqrt{n^2 + m^2}$$

$$\text{Need } \begin{cases} u(x, y, 0) = \sum_{n, m=1}^{\infty} A_{nm} \sin nx \sin my \quad \leftarrow \text{Get A's} \end{cases}$$

$$\begin{cases} \frac{\partial u}{\partial t}(x, y, 0) = \sum_{n, m=1}^{\infty} \omega_{nm} B_{nm} \sin nx \sin my \quad \leftarrow \text{Get B's} \end{cases}$$

$$\text{Need } \sum_{n=1}^{\infty} \left(\underbrace{\sum_{m=1}^{\infty} A_{nm} \sin my}_{\substack{\text{Fourier sine series} \\ \text{coeff for } f(\cdot, y)}} \right) \sin nx = f(x, y) \quad \text{want}$$

$$\text{See } \left(\sum_{m=1}^{\infty} A_{nm} \sin my \right) = \underbrace{\frac{2}{\pi} \int_0^{\pi} f(x, y) \sin nx \, dx}_{= F(y)}$$

$$A_{nm} = \frac{2}{\pi} \int_0^{\pi} F(y) \sin my \, dy$$

$$A_{nm} = \frac{2}{\pi} \int_0^{\pi} \left(\frac{2}{\pi} \int_0^{\pi} f(x, y) \sin nx \, dx \right) \sin my \, dy = \quad \text{Fubini's}$$

$$A_{nm} = \frac{4}{\pi^2} \int_0^{\pi} \int_0^{\pi} f(x,y) \sin nx \sin my \, dx \, dy$$

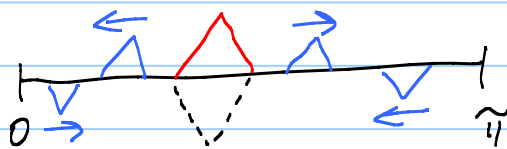
Similarly, $B_{nm} = \frac{4}{\pi^2 \omega_{nm}} \int_0^{\pi} \int_0^{\pi} g(x,y) \sin nx \sin my \, dx \, dy$

Alternatively: $\sin nx \sin my$ are orthogonal:
 $\mathcal{Q}_{nm}$

$$\begin{aligned} \langle \mathcal{Q}_{nm}, \mathcal{Q}_{kl} \rangle &= \int_0^{\pi} \int_0^{\pi} \mathcal{Q}_{nm}(x,y) \mathcal{Q}_{kl}(x,y) \, dx \, dy \\ &= \begin{cases} 0 & (n,m) \neq (k,l) \\ \neq 0 & (n,m) = (k,l) \end{cases} \end{aligned}$$

Fun facts: $\omega_{nm} = \sqrt{n^2 + m^2}$ discordant overtones!

Another strange thing: String: $\omega_n = n \leftarrow \cos nt$ 2π -periodic
 Sq. drum: ω_{nm} don't have a common period



Initial shape gets repeated!



infinite string

Huygen's principle: waves in $\mathbb{R}^1, \mathbb{R}^3, \mathbb{R}^5, \dots$ go away, leave space in peace
 $\mathbb{R}^2, \mathbb{R}^4, \mathbb{R}^6, \dots$ they don't

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Laplace transform: $f(x) \mapsto F(s) = \int_0^{\infty} f(x) e^{-sx} dx$

Useful for studying initial value problems at $x=0$ for ODEs.

Turns linear ode's with const coeff into algebra!

Big facts: $\mathcal{L}[f(x)] = F(s)$ is a linear operator.

1) $\mathcal{L}[f'] = sF(s) - f(0)$

2) $\mathcal{L}[f''] = s^2 F(s) - sf(0) - f'(0)$

3) Can undo the Laplace transform!

If $\mathcal{L}[f_1] = \mathcal{L}[f_2]$ for all $s > M$ (for some M),

then $f_1 \equiv f_2$ on $[0, \infty)$.

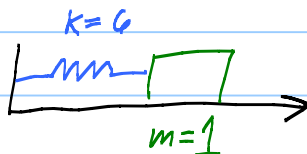
Consequently, $\mathcal{L}^{-1}$ makes sense.

Plan: $\mathcal{L}[\text{ODE}] = \text{Algebra in } F(s)$

Solve for F . Solⁿ: $f(x) = \mathcal{L}^{-1}[F(s)]$.

EX: $y'' + 5y' + 6y = e^{-4t}$

$\uparrow$
 $c=5$ friction
term



$$m y'' = -\underset{\substack{\uparrow \\ \text{Hooke's}}}{6} y - \underset{\substack{\uparrow \\ \text{friction}}}{5} y' + \underset{\substack{\uparrow \\ \text{external}}}{e^{-4t}}$$

$$\mathcal{L}[y'' + 5y' + 6y] = \mathcal{L}[e^{-4t}]$$

$$(s^2 \bar{Y}(s) - sy(0) - y'(0)) + 5(s \bar{Y}(s) - y(0)) + 6 \bar{Y} = \frac{1}{s+4}$$

$$\underbrace{(s^2 + 5s + 6)}_{(s+2)(s+3)} \bar{Y}(s) = sy(0) + y'(0) + 5y(0) + \frac{1}{s+4}$$

Plug in initial conditions. Solve for $\bar{Y}$.

$$\text{Take } y(0)=0, y'(0)=0$$

$$\text{Get } \bar{Y}(s) = \frac{1}{(s+2)(s+3)(s+4)}$$

$$\stackrel{\text{P.F.}}{=} \frac{A}{s+2} + \frac{B}{s+3} + \frac{C}{s+4}$$

$$y(t) = Ae^{-2t} + Be^{-3t} + Ce^{-4t}$$