

MA 428 Exam 2 solutions

$$\begin{aligned} 1. \quad \hat{f}(s) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-isx} dx \\ &= \frac{1}{\sqrt{2\pi}} \left(\int_0^1 \underbrace{x}_{u} \underbrace{e^{-isx}}_{dv} dx + \int_1^2 \underbrace{(2-x)}_u \underbrace{e^{-isx}}_{dv} dx \right) \\ &= \frac{1}{\sqrt{2\pi}} \left[x \cdot \left(\frac{-1}{is} e^{-isx} \right) \right]_0^1 - \int_0^1 \left(\frac{-1}{is} e^{-isx} \right) dx \\ &\quad + (2-x) \left(\frac{-1}{is} e^{-isx} \right) \Big|_1^2 - \int_1^2 \left(\frac{-1}{is} e^{-isx} \right) (-dx) \Big] \\ &= \frac{1}{\sqrt{2\pi}} \left(\frac{-e^{-is}}{is} - \left[\frac{1}{(is)^2} e^{-isx} \right]_0^1 \right. \\ &\quad \left. - \left(\frac{-e^{-is}}{is} \right) + \left[\frac{1}{(is)^2} e^{-isx} \right]_1^2 \right) \\ &= \frac{1}{\sqrt{2\pi}} \left[\frac{1}{s^2} (e^{-is} - 1) - \frac{1}{s^2} (e^{-i2s} - e^{-is}) \right] = \end{aligned}$$

$$\hat{f}(s) = \frac{1}{\sqrt{2\pi}} \left(\frac{-e^{-i2s} + 2e^{-is} - 1}{s^2} \right)$$

Since $\sqrt{2\pi} \hat{f}(s) = \int_{-\infty}^{\infty} = \int_0^{\infty} f(x) \underset{\substack{\uparrow \\ \cos(sx) - i\sin(sx)}}{e^{-isx}} dx,$

we see that $\mathcal{N}_c f = \sqrt{\frac{2}{\pi}} \sqrt{2\pi} \operatorname{Re} \hat{f}(s).$

$$\begin{aligned} \text{So } [\mathcal{N}_c f](s) &= \sqrt{\frac{2}{\pi}} \operatorname{Re} \left(\frac{-e^{-i2s} + 2e^{-is} - 1}{s^2} \right) \\ &= \sqrt{\frac{2}{\pi}} \left(\frac{-\cos 2s + 2\cos s - 1}{s^2} \right) \end{aligned}$$

2. Let $\hat{u}(s, t) = \mathcal{N}_c [u(x, t)](s)$

$$= \frac{1}{\sqrt{2\pi}} \int_0^{\infty} u(x, t) \cos(sx) dx. \quad \text{Then}$$

$$\frac{\partial \hat{u}}{\partial t} = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \underbrace{\frac{\partial u}{\partial t}(x, t)}_{= \frac{\partial^2 u}{\partial x^2}} \cos(sx) dx$$

$$= \mathcal{F}_c \left[\frac{\partial^2 u}{\partial x^2} \right] = -s^2 \underbrace{\mathcal{F}_c[u]}_{\hat{u}} - \underbrace{\sqrt{\frac{2}{\pi}} \frac{\partial u}{\partial x}(0, t)}_{=0}$$

So $\boxed{\frac{\partial \hat{u}}{\partial t} = -s^2 \hat{u}}$ and $\hat{u}(s, t) = C(s) e^{-s^2 t}$.

Plug in $t=0$ to see $C(s) = \hat{u}(s, 0) = [\mathcal{F}_c f](s) =$

$$\sqrt{\frac{2}{\pi}} \left(\frac{-\cos 2s + 2\cos s - 1}{s^2} \right). \quad \text{Finally,}$$

$$u(x, t) = \mathcal{F}_c^{-1} \hat{u} = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \sqrt{\frac{2}{\pi}} \left(\frac{-\cos 2s + 2\cos s - 1}{s^2} \right) e^{-s^2 t} \cos(sx) ds$$

3. Let $\hat{u}(s, t) = \frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} u(x, t) e^{-isx} dx$.

$$\frac{\partial \hat{u}}{\partial t} = \frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} \underbrace{\frac{\partial u}{\partial t}(x, t)}_{= \frac{\partial^2 u}{\partial x^2}} e^{-isx} dx$$

$$= \mathcal{F} \left[\frac{\partial^2 u}{\partial x^2} \right] = -s^2 \mathcal{F}[u]$$

So $\frac{\partial \hat{u}}{\partial t} = -s^2 \hat{u}$ and $\hat{u}(s, t) = C(s) e^{-s^2 t}$

where, this time, $C(s)$ needs to be an arbitrary complex valued function of s .

Plug in $t=0$ to see $C(s) = \hat{u}(s, 0) =$

$$\frac{1}{\sqrt{2\pi i}} \int_{-\infty}^{\infty} \underbrace{u(x, 0)}_{f(x)} e^{-isx} dx = \hat{f}(s). \text{ Finally,}$$

$$u(x,t) = \mathcal{N}^{-1} \left[\hat{f}(s) e^{-s^2 t} \right]$$

$$\underline{\text{ans}} \rightarrow = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi}} \left(\frac{-e^{-i2s} + 2e^{-is} - 1}{s^2} \right) e^{-s^2 t + isx} ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\underbrace{\frac{-\cos 2s + 2\cos s - 1}{s^2}}_{\text{even}} + i \underbrace{\frac{\sin 2s - 2\sin s}{s^2}}_{\text{odd}} \right] \underbrace{e^{-s^2 t}}_{\text{even}} ds$$

$$\cdot \left[\underbrace{\cos sx}_{\text{even}} + i \underbrace{\sin sx}_{\text{odd}} \right] ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\frac{-\cos 2s + 2\cos s - 1}{s^2} \cos(sx) - \frac{\sin 2s - 2\sin s}{s^2} \sin(sx) \right] e^{-s^2 t} ds$$

4. D'Alembert's method calls for making the change of variables $w = x + ct$, $v = x - ct$ to transform the wave equation $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ to the easy equation $\frac{\partial^2 u}{\partial w \partial v} = 0$, which has solution $u(w, v) = \Phi(w) + \Psi(v)$, yielding

$$u(x, t) = \Phi(x + ct) + \Psi(x - ct).$$

Note that $\frac{\partial u}{\partial t}(x, t) = c\Phi'(x + ct) - c\Psi'(x - ct)$,

so, to satisfy the initial conditions

$$\begin{cases} u(x, 0) = \Phi(x) + \Psi(x) \stackrel{\text{want}}{=} e^{-x^2} & (A) \\ \frac{\partial u}{\partial t}(x, 0) = c\Phi'(x) - c\Psi'(x) = 0 & (B) \end{cases}$$

We need eqns (A) and (B) to hold.

(B) shows that $\phi' \equiv \psi'$, so

$$\psi(x) = \phi(x) + K, \text{ where } K = \text{const.}$$

Now (A) yields $\phi(x) + [\phi(x) + K] = e^{-x^2}$,

which we can solve by setting $K=0$ and $\phi(x) = \frac{1}{2}e^{-x^2}$. Note, since $K=0$,

we have $\psi(x) = \phi(x) = \frac{1}{2}e^{-x^2}$ too.

$$\text{Finally, } u(x,t) = \frac{1}{2} \left[e^{-(x+ct)^2} + e^{-(x-ct)^2} \right]$$

which represents the hump e^{-x^2} breaking into two half size humps traveling in $\pm$ directions at speed c .

Remarks $\hat{f}(s)$ and $\mathcal{N}_{\mathbb{C}}[f]$ are not singular at $s=0$ because, if you write out

$$e^{-i2s} = 1 + (-i2s) + \frac{(-i2s)^2}{2!} + \dots$$

$$e^{-is} = 1 + (-is) + \frac{(-is)^2}{2!} + \dots$$

and collect terms in

$$-e^{-i2s} + 2e^{-is} - 1 = 0 + 0 \cdot s + 1 \cdot s^2 + \dots$$

$$= s^2 \cdot (\text{convergent power series on } \mathbb{R}),$$

you can deduce that $\hat{f}(s)$ and $\mathcal{N}_{\mathbb{C}}[f]$ are C^∞ -smooth on $\mathbb{R}$ when given the proper values at $s=0$.